

New Phenomenology for Palatini f(R): Non-singular Universes

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About this talk ...

Modified theories of gravity of the *f*(*R*) type have been thoroughly studied in the recent literature in connection with the cosmic speedup problem.

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- Modified theories of gravity of the f(R) type have been thoroughly studied in the recent literature in connection with the cosmic speedup problem.
- I will show here that Palatini f(R) theories can also be used to address issues of the very early Universe such as the Big Bang singularity.

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- Structure of the talk:
 - Basics of the Palatini f(R) field equations.
 - What is new in these theories?
 - Infrared and Ultraviolet corrected models.
 - Applications to the early Universe and Quantum Gravity.



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Basics of Palatini and IR-Corrected Models



• The action is defined as
$$S[g,\Gamma,\psi_m] = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} f(R) + S_m[g_{\mu\nu},\psi_m]$$

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- The action is defined as $S[g,\Gamma,\psi_m] = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} f(R) + S_m[g_{\mu\nu},\psi_m]$
- In Palatini, metric and connection are independent fields:
 - Metric variation: $f_R R_{\mu\nu}(\Gamma) \frac{f}{2}g_{\mu\nu} = \kappa^2 T_{\mu\nu}$ where $f_R = df/dR$ Connection variation: $\tilde{\nabla}_{\beta}[\sqrt{-g}f_R g^{\mu\nu}] = 0$

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- Using the trace $Rf_R 2f = \kappa^2 T$, which implies $R = \mathcal{R}(T)$, it follows that $\Gamma^{\alpha}_{\mu\nu}$ is the Levi-Civita connection of $h_{\mu\nu} = f_R(T)g_{\mu\nu}$.

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- The metric equations can be written in two alternative ways. Using $g_{\mu\nu}$:

$$G_{\mu\nu}(g) = \frac{\kappa^2}{f_R} T_{\mu\nu} - \frac{Rf_R - f}{2f_R} g_{\mu\nu} + \frac{1}{f_R} \left(\nabla_\mu \nabla_\nu f_R - g_{\mu\nu} \Box f_R \right) - \frac{3}{2f_R^2} \left(\partial_\mu f_R \partial_\nu f_R - \frac{1}{2} g_{\mu\nu} (\partial f_R)^2 \right)$$

Or using $h_{\mu\nu}$: $G_{\mu}{}^{\nu}(h) = \frac{\kappa^2}{f_R^2(T)} \left[T_{\mu}{}^{\nu} - \left(\frac{Rf_R - f}{2\kappa^2} \right) \delta_{\mu}{}^{\nu} \right]$

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In vacuum, both reduce to $G_{\mu\nu} = -\Lambda_{eff}g_{\mu\nu}$,

where
$$\Lambda_{eff} = \frac{Rf_R - f}{2f_R}\Big|_0$$

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- Despite f(R) being non-linear, we still have second-order equations.
- Birkhoff's theorem holds (same vacuum solutions as GR)

$$ds^{2} = -\left(1 - \frac{2M}{r} - \frac{\Lambda_{eff}r^{2}}{3}\right)dt^{2} + \frac{dr^{2}}{\left(1 - \frac{2M}{r} - \frac{\Lambda_{eff}r^{2}}{3}\right)} + r^{2}d\Omega^{2}$$

Scharzschild-de Sitter is the most general spherical, vacuum solution.

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 $g_{\mu\nu} \text{ has two types of contributions: from the total amount of matter and from the local energy-momentum densities: } g_{\mu\nu} = \frac{1}{f_R(T)} h_{\mu\nu} .$

Note that $G_{\mu}^{\nu}(h) = \frac{\kappa^2}{f_R^2(T)} \left[T_{\mu}^{\nu} - \left(\frac{Rf_R - f}{2\kappa^2} \right) \delta_{\mu}^{\nu} \right]$ has the same differential structure as GR.

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IR-corrected models are problematic: atoms are unstable ([G.J.O., PRD77(2008)]) In $f(R) = R - \frac{\mu^4}{R}$ Hydrogen would disintegrate in less than 2 hours !!!

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- IR-corrected models are problematic: atoms are unstable ([G.J.O., PRD77(2008)]) In $f(R) = R - \frac{\mu^4}{R}$ Hydrogen would disintegrate in less than 2 hours !!!
- UV-corrected models may be designed to avoid the Big Bang singularity.
 - Bouncing and cyclic cosmologies are very common in Palatini f(R).
 - There exist important links between Palatini and the effective dynamics of Loop Quantum Cosmology.

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• For very weak sources, $ds^2 \approx \frac{1}{f_R(T)} \eta_{\mu\nu} dx^{\mu} dx^{\nu}$.

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• For very weak sources,
$$ds^2 \approx \frac{1}{f_R(T)} \eta_{\mu\nu} dx^{\mu} dx^{\nu}$$

The curved-space Dirac equation $(i\lambda^{\mu}D_{\mu} - m)\psi = 0$ turns into

$$i\partial_t \Psi = \left[\vec{\alpha} \cdot (\vec{p} - e\vec{A} + i\vec{\nabla}\Omega) + (eA_0 + \partial_t\Omega) + \tilde{m}\beta\right]\Psi$$

where
$$\Omega \equiv (3/4) \ln \phi(T)$$
, $\tilde{m} \equiv m \phi^{-\frac{1}{2}}$, $\phi \equiv f_R(T)/f_R(0)$, and $T = -m \bar{\psi} \psi$

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For stationary states $\psi = e^{-iEt}\xi(\vec{x})$, we have $E\xi = [\vec{\alpha} \cdot \vec{\pi} + eA_0 + \tilde{m}\beta]\xi$.

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• The positive energy spinor η in the non-relativistic limit satisfies:

$$\mathcal{E}\eta = \left\{ \frac{1}{\tilde{m}+m_0} \left[(\vec{p} - e\vec{A})^2 - e\vec{\sigma} \cdot \vec{B} \right] + eA_0 + V(\Omega) + \Delta m \right\} \eta$$

where
$$E = m_0 + \mathcal{E}$$
, $T \approx -m\eta^{\dagger}\eta$, $\Delta m = \tilde{m} - m_0$, and

$$V(\Omega) = \frac{1}{\tilde{m} + m_0} \left[i \vec{\sigma} (\vec{\nabla} \Omega \times \vec{\nabla}) - 2ie(\vec{A} \cdot \vec{\nabla} \Omega) + 2(\vec{\nabla} \Omega \cdot \vec{\nabla}) + \vec{\nabla}^2 \Omega - |\vec{\nabla} \Omega|^2 \right]$$

The gravity lagrangian induces new interaction terms $V(\Omega)$ and Δm

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Ground state $\eta_{(1,0,0)}$ for $f(R) = R - \mu^4/R$



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Ground state $\eta_{(1,0,0)}$



- The potential well induced by Δm tends to $\Delta m \approx -0.13m \approx -5000 \times 13'6 \text{ eV}$
- This well triggers the flux of probability density to the exterior of the atom.
- The electron escapes similarly as in nuclear alpha decay.
- In the 1/R model, the half life of the electron is $\tau_e \sim 6 \times 10^3$ s.

Non-perturbative effects in $\eta_{(2,0,0)}$



Non-perturbative effects in $\eta_{(2,0,0)}$

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IR-corrected models are characterized by a low density scale which is reached at late times in cosmological models, when the averaged energy density is very low.

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- Such energy scales can also be reached in atomic systems like Hydrogen.

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- Strong non-perturbative gravitational effects arise near the zeros (radial and angular!) of the wave functions.

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Strong gravitational effects occur in IR models when their characteristic density scale is reached. Are there similar effects in UV models?



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• Characterizing the Bounce

• Palatini f(R) and LQC

• Numerics and Fits

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UV models and Non-singular Universes

• In the simplest UV model,
$$f(R) = R \pm \frac{R^2}{R_P}$$
, we find

$$H^2 = \frac{\kappa^2 \rho}{3} \frac{\left(1 + \frac{2(1-3w)\kappa^2 \rho}{R_P}\right) \left(1 + \frac{(1-3w)^2}{2} \frac{\kappa^2 \rho}{R_P}\right)}{\left[1 - (1-9w^2)\frac{\kappa^2 \rho}{R_P}\right]^2} - \frac{K}{a^2}$$

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Non-singular Universes may arise depending on the EoS $P = w\rho$



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For a general f(R) theory, the Hubble function is given by ($P = w\rho$)

$$3H^2 = \frac{f_R\left(\kappa^2\rho + \frac{Rf_R - f}{2}\right)}{\left(f_R + \frac{3}{2}\Lambda_1\right)^2} - \frac{3K}{a^2}$$

where
$$\Lambda_1 = (1+w)(1-3w)\kappa^2 \rho \frac{f_{RR}}{Rf_{RR}-f_R}$$

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Cosmic bounce occurs whenever $H^2 = 0$. When K = 0, this may happen if [C.Barragán, G.J.O, and H.Sanchis-Alepuz, To appear]

- First case: $f_R \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.
- Second case: $\kappa^2 \rho + \frac{Rf_R f}{2} \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.

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- First case: $f_R \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.
- Second case: $\kappa^2 \rho + \frac{Rf_R f}{2} \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.
- When $K \neq 0$, the discussion is not as clean and requires a model by model study.

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For a general f(R) theory, the Hubble function is given by ($P = w\rho$)

$$3H^2 = \frac{f_R\left(\kappa^2\rho + \frac{Rf_R - f}{2}\right)}{\left(f_R + \frac{3}{2}\Lambda_1\right)^2} - \frac{3K}{a^2} \quad \text{where} \quad \Lambda_1 = (1+w)(1-3w)\kappa^2\rho \frac{f_{RR}}{Rf_{RR} - f_R}$$

Cosmic bounce occurs whenever $H^2 = 0$. When K = 0, this may happen if [C.Barragán, G.J.O, and H.Sanchis-Alepuz, To appear]

- First case: $f_R \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.
- Second case: $\kappa^2 \rho + \frac{Rf_R f}{2} \to 0$ and $f_R + \frac{3}{2}\Lambda_1 > 0$.
- When $K \neq 0$, the discussion is not as clean and requires a model by model study.

It is possible to construct Palatini f(R) theories as successful as GR at low energies but free from high curvature singularities. Is this just an accident or is there anything deeper?

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- A Non-singular Universe
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- Palatini f(R) and LQC
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- Summary and Conclusions



Loop Quantum Cosmology is an approach based on the quantization techniques of canonical quantum gravity (LQG).

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Loop Quantum Cosmology is an approach based on the quantization techniques of canonical quantum gravity (LQG).

Though the fundamental equations are discrete, in the continuum limit one can find the following effective o.d.e. :

$$H^2 = \frac{8\pi G}{3} \rho \left(1 - \frac{\rho}{\rho_{crit}} \right)$$
, with $\rho_{crit} = 0.41 \rho_{Planck}$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}\rho\left(1 - 4\frac{\rho}{\rho_{crit}}\right) - 4\pi GP\left(1 - 2\frac{\rho}{\rho_{crit}}\right)$$

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No new degrees of freedom in LQC. Matter alone can cure the singularity.

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- Though the fundamental equations are discrete, in the continuum limit one can find the following effective o.d.e. :

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- No new degrees of freedom in LQC. Matter alone can cure the singularity.
- The non-perturbative dynamics of LQC can be captured by a Palatini f(R).
 - Use the trace equation $(3w-1)\kappa^2\rho = Rf_R 2f$ to find $\rho = \rho(R, f, f_R)$.
 - For K = 0 and w = 1: $8\pi G\rho\left(1 \frac{\rho}{\rho_{crit}}\right) = \frac{f_R(\kappa^2 \rho + (\mathcal{R} f_R f)/2)}{\left(f_R \frac{12\kappa^2 \rho f_{RR}}{2(\mathcal{R} f_{RR} f_R)}\right)^2}$.

• Simple manipulations lead to the following o.d.e.:

$$f_{RR} = -f_R \left(\frac{Af_R - B}{2(\mathcal{R} f_R - 3f)A + \mathcal{R} B} \right)$$

where

$$A = \sqrt{2(\mathcal{R} f_R - 2f)(2\mathcal{R}_c - [\mathcal{R} f_R - 2f])},$$

$$B = 2\sqrt{\mathcal{R}_c f_R(2\mathcal{R}_c f_R - 3f)}, \text{ and } \mathcal{R}_c \equiv \kappa^2 \rho_c.$$

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Dashed red line: Numerical Curve. [G.J.O and P.Singh, JCAP030(2009)]



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Dashed red line: Numerical Curve. [G.J.O and P.Singh, JCAP030(2009)]



This f(R) lagrangian exactly reproduces the dynamics of isotropic LQC.

The cosmic bounce occurs at $R = -12R_c$, where $f_R \to 0$.

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This f(R) lagrangian exactly reproduces the dynamics of isotropic LQC.

The cosmic bounce occurs at $R = -12R_c$, where $f_R \to 0$.

Palatini f(R) is able to fully capture the dynamics of LQC even though the form of the lagrangian near the bounce cannot be estimated by perturbative methods.

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Palatini f(R) theories modify the gravitational dynamics without adding any new degrees of freedom.

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IR and **UV** corrections w.r.t. the E-H lagrangian remain (almost) invisible until their respective density scales are reached. The new effects appear non-perturbatively.

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- LQC can be exactly reproduced by a Palatini f(R) action. Palatini theories could play an important role in phenomenological descriptions of quantum gravity.

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- LQC can be exactly reproduced by a Palatini f(R) action. Palatini theories could play an important role in phenomenological descriptions of quantum gravity.
- Future work:
 - Are black holes non-singular in Palatini f(R)?
 - How is the dynamics of extended Palatini theories $f(R, R_{\mu\nu}R^{\mu\nu}, ...)$?



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Thanks !!!

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