

June 29 - July 3, 2009

NATURAL NEUTRINO DARK ENERGY Ilya Gurwich Physics Department Ben-Gurion University

Wednesday, July 1, 2009

Outline

- * Motivation
- * Some general results for a new class of models
- * Results vs measurements
- * An example model
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- * Open question and future research

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Despite improving observational constraints, constantly used to sort the different models, "theoreticians work faster".

It is important to see whether we can solve dark energy without invoking new physics and new particles. And what is the most we can say without constraining ourselves to a specific model.

Motivation - Energy Scales

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* $\Lambda \sim (10^{-33} eV)^2$: No known scale to match.

* $\rho_{\Lambda} \sim (10^{-3} eV)^4$: The lightest neutrino mass, m_{ν} , matches the scale. **The only fundamental scale** that matches.

* There are other complex/composite (non-fundamental) scales that match this value, for example: $\Lambda_{QCD}^2/M_{plank}$

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> All models suffer from one common problem What is so special about the neutrino?

- * The lightest neutrino is (probably) the only massive particles that is still relativistic... So what?
- * non-relativistic particles have a relatively constant energy \simeq m.
- * relativistic particles have an energy that changes as the universe expands $\sim a(t)^{-1}$.

- * From a QFT point of view, one may view the particles as spheres of size E^{-1} , where E is the particle energy.
- * For relativistic particles, the concentrations of these spheres $\sim a^3 N(E^{-1})^3 \propto (El)^{-3}$ (N is the number of particles and l is the mean distance between them), meaning the relative space they occupy is constant as the universe expands.
- * It is also important to note that the quantity El, governs the propagation of the particles over the average distance between them l, so it is expected to govern any relevant interaction that is expected to lead to the dark energy contribution.

a(t')=2a(t)

For relativistic particles, the situation stays constant in the comoving frame.

For non-relativistic particles, from the comoving frame, it seems as if they disappear, any interaction between them would weaken.

ultra-relativistic

non-relativistic

a(t)

- * Another relevant parameter that behaves as a constant for relativistic particles and grows for non-relativistic ones is E/T.
- * This factor may become relevant for creation and annihilation operators of states that generate the dark energy, as $\exp(-E/T)$.

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$$\rho_{DE} = \sum_{fermions} f(\xi_i)$$

Where: $\xi_i = E_i/T_i$ or $\xi_i = E_i l_i$

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For the dark energy to be naturally related to relativistic particles and fit the mass scale, we also assume:

$$\lim_{x \to 0} f(x) \sim m^4 \quad \text{and} \quad \lim_{x \to \infty} f(x) = 0$$

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We have:

$$1 + w_{DE} = -\frac{1}{3} \frac{a}{\rho_{DE}} \frac{\partial \rho_{DE}}{\partial a} = -\frac{1}{3} \frac{a}{\rho_{DE}} f'(E_{\nu}l_{\nu}) \frac{\partial(E_{\nu}l_{\nu})}{\partial a}$$

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Eventually we have: $E_{\nu}l_{\nu} \simeq \alpha + \beta \left(\frac{m_{\nu}}{T_{\nu}}\right)^2$

$$\alpha = \frac{7}{90} \left(\frac{\pi^{14}}{12\zeta(3)^4} \right)^{1/3} \simeq 5.554 \qquad \beta = \frac{1}{18} \left(\frac{\pi^8}{12\zeta(3)^4} \right)^{1/3} \simeq 0.402$$

$$1 + w_{DE} \simeq -\frac{2\beta}{3} \frac{f'(E_{\nu}l_{\nu})}{f(E_{\nu}l_{\nu})} \left(\frac{m_{\nu}}{T_{\nu}}\right)^{2}$$

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Eventually we have: $E_{\nu}/T_{\nu} \simeq \alpha + \beta \left(\frac{m_{\nu}}{T_{\nu}}\right)^2$

$$\alpha = \frac{7\pi^4}{180\zeta(3)} \simeq 3.151 \qquad \beta = \frac{5\pi^2}{180\zeta(3)} \simeq 0.228$$

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Natural Neutrino Dark Energy - Results vs Observations

Conventional notations: $w_a = -2(1 + w_0)$



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Although there is no clear conclusion from todays data - there is room for optimism!

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Natural Neutrino Dark Energy vs Other models



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Natural Neutrino Dark Energy vs Quintessence

from R. Caldwell and M. Kamionkowski (2009)



Compared to thawing and freezing quintessence scenarios, the neutrino dark energy is compatible with the thawing scenario. So how can we distinguish?

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For present data this may represent too many parameters So a bias parameterization may be

$$w = w_0 + w_1 \left[(1-a) - \frac{1}{2}(1-a)^2 \right]$$

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Which happens first?



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Natural Neutrino Dark Energy Cosmology



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$$1 + w_{DE} \simeq \left(\frac{m_{\nu}}{3.76 \cdot 10^{-4}eV}\right)^{2}$$

Coupled vs Decoupled

Decoupled (assuming decoupling temperature >> m)	Coupled
$n = 2 \int_0^\infty \frac{4\pi p^2 dp}{(2\pi)^3} \frac{1}{1 + \exp(p/T)}$	$n = 2 \int_0^\infty \frac{4\pi p^2 dp}{(2\pi)^3} \frac{1}{1 + \exp(\sqrt{p^2 + m^2}/T)}$
$\rho = 2 \int_0^\infty \frac{4\pi p^2 dp}{(2\pi)^3} \frac{\sqrt{p^2 + m^2}}{1 + \exp(p/T)}$	$\rho = 2 \int_0^\infty \frac{4\pi p^2 dp}{(2\pi)^3} \frac{\sqrt{p^2 + m^2}}{1 + \exp(\sqrt{p^2 + m^2}/T)}$
	coupled fermions decoupled fermions

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The neutrino and quark masses used in the plot are speculative

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Fermion mass hierarchy generates a long period of dark energy domination. However because some of the fermion masses are of the same order of magnitude, DE domination does not imply necessarily that $w_{eff} \simeq -1$.



For the general class of models

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- * Since dark energy domination leads to inflationary expansion, it may have an effect on inflationary constraints. Can such an effect be measured?
- * Will tighter dark energy constraints from future measurements, verify that $w_a = -2(1 + w_0)$?

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- * Precise measurement of neutrino masses using cosmology.
- * Does neutrino non-homogeneity (clustering) effects the dark energy? Does it lead to non-homogeneity in dark energy?
- * Can the interactions that lead to dark energy be measured in the laboratory?

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- * A model following the above guidelines was presented.
- * Many intriguing questions for future research follow from this work.


THANK YOU FOR YOUR ATTENTION